

ARTICLE



<https://doi.org/10.1038/s41467-021-21831-y>

OPEN

Fragility of surface states in topological superfluid ^3He

P. J. Heikkinen¹✉, A. Casey¹, L. V. Levitin¹, X. Rojas¹, A. Vorontsov², P. Sharma³, N. Zhelev⁴, J. M. Parpia⁴ & J. Saunders¹✉

Superfluid ^3He , with unconventional spin-triplet p-wave pairing, provides a model system for topological superconductors, which have attracted significant interest through potential applications in topologically protected quantum computing. In topological insulators and quantum Hall systems, the surface/edge states, arising from bulk-surface correspondence and the momentum space topology of the band structure, are robust. Here we demonstrate that in topological superfluids and superconductors the surface Andreev bound states, which depend on the momentum space topology of the emergent order parameter, are fragile with respect to the details of surface scattering. We confine superfluid ^3He within a cavity of height D comparable to the Cooper pair diameter ξ_0 . We precisely determine the superfluid transition temperature T_c and the suppression of the superfluid energy gap, for different scattering conditions tuned in situ, and compare to the predictions of quasiclassical theory. We discover that surface magnetic scattering leads to unexpectedly large suppression of T_c , corresponding to an increased density of low energy bound states.

¹Department of Physics, Royal Holloway University of London, Surrey, UK. ²Department of Physics, Montana State University, Bozeman, MT, USA.

³Department of Physics, Indian Institute of Science, Bangalore, India. ⁴Department of Physics, Cornell University, Ithaca, NY, USA. ✉email: petri.heikkinen@rhul.ac.uk; j.saunders@rhul.ac.uk

The spin-triplet superfluid phases^{1–3} of liquid ³He provide a benchmark for topological superconductivity⁴. So far such superconductors^{5,6} are the missing ‘elements’ in the periodic table of quantum matter⁷; while candidates exist, such as Sr₂RuO₄⁸, UPT₃⁹, doped Bi₂Se₃¹⁰, UTe₂¹¹, no bulk material has yet been unambiguously identified as a topological crystalline superconductor¹². Nevertheless, devices fabricated from spin-triplet crystalline topological superconductors should eventually contribute to potential applications in topological quantum computation^{13–15}. Current strategies explore the manipulation of Majorana zero modes (MZMs) in devices which rely on 1-D topological superconductivity^{16–19} induced by the proximity effect in topological insulators or semiconductors with strong spin–orbit coupling, or via planar Josephson Junctions^{20,21}. However, MZMs may be corrupted by non-topological low energy Andreev bound states (ABS)²², which can be present as a result of details of interfaces and materials properties in such systems²³. Here we report the fragility of surface ABS in superfluid ³He at an ideal non-transparent interface, exploiting the ability to tune in situ the scattering of quasiparticles by the surface through adjustment of the isotopic composition of the helium surface boundary layer.

Recently we have shown that it is possible to cool ³He confined within precisely engineered nanoscale cavities into the superfluid phases²⁴, and detect the nuclear magnetic resonance (NMR) response of the small ³He sample using an ultra-sensitive spectrometer²⁵. Surface scattering dominates the properties under strong confinement. NMR determines the superfluid transition temperature, the pairing state, and the superfluid energy gap^{24,26,27}, the suppression of which under confinement self-consistently determines, through quasiclassical theory^{28,29}, the density of states of the ABS mid-gap surface excitations.

Superfluid ³He consists of spin-triplet Cooper pairs, with one unit ($l=1$) of orbital angular momentum¹. The order parameter is a complex 3×3 matrix, encoding the spin state of the pairs over the spherical Fermi surface. In bulk liquid, two phases are found with distinct symmetries^{1,2} and momentum-space topologies⁴. The A phase is chiral, breaking time-reversal symmetry. Over the Fermi surface, pairs form with the same direction of their orbital angular momentum and in an equal-spin state comprising just $|\uparrow\uparrow\rangle$ and $|\downarrow\downarrow\rangle$ pairs. The B phase is time-reversal invariant, comprising all three components of the spin triplet, with broken relative spin–orbit symmetry. The relative stability of these phases is profoundly altered by confinement²⁴. Superfluid ³He is an intrinsically impurity-free system, although impurities can be artificially introduced using silica aerogels of different porosities and structure factors³⁰. Our study of superfluid ³He confined in a simple slab geometry determines the influence of surface scattering alone on gap suppression, in the absence of impurity scattering.

In the quasiclassical theory of superfluid ³He²⁹ the effect of the surface can be characterized in terms of a single parameter, S , the degree of specularity of the surface scattering of quasiparticles³¹. There is prior compelling evidence from hydrodynamic and transverse acoustic impedance studies of normal state ³He that the surface scattering may be tuned in situ from diffuse to specular^{32–34} by coating surfaces with ⁴He. Approaching close to full specularity was found to require superfluidity of the ⁴He surface film^{33,35}. In early work it was shown that both the NMR frequency shift of superfluid ³He confined in a stack of mylar sheets³⁶ and the superfluid fraction within the pores of packed powders³⁷ depend on the surface ⁴He coverage. They both increased with increased coverage, indicating the expected reduced pair breaking with higher specularity. Transverse acoustic impedance measurements, and quasiclassical analysis, revealed in superfluid ³He-B the energy density of states of mid-

gap surface-bound excitations depends on specularity^{38–41}. Their spectrum shows a Majorana-like cone as the conditions for fully specular scattering are approached.

Quasiclassical theory predicts for the A phase that T_c and gap suppression, and surface-bound states are all eliminated for fully specular scattering. Diffuse scattering leads to a finite density of surface-bound states at low energy. In this case, the suppression of T_c scales as $\delta T_c/T_{c0} \propto -(\xi_0/D)^2$ to leading order. In the work reported here, we use NMR to study the A phase in a single cavity of precisely defined dimensions, with several different surface boundary layers to tune the surface quasiparticle scattering.

Results

Experimental details. In our experiment, superfluid ³He was confined within a 192 nm high cavity defined in a silicon wafer, Fig. 1a. The effective confinement can be varied at fixed cavity height D by changing the sample pressure and hence the superfluid coherence length $\xi_0 = \hbar v_F / 2\pi k_B T_{c0}$, where v_F is the Fermi velocity and T_{c0} the bulk superfluid transition temperature. Measurements were made at a series of pressures from 0.0 to 5.5 bar, over which ξ_0 decreases from 77 to 40 nm. We determine the shift in the NMR resonance frequency f relative to the Larmor frequency f_L , $\Delta f = f - f_L$, which occurs in the superfluid state. The onset of this shift identifies T_c in the cavity. This is determined precisely relative to T_{c0} by also observing frequency shifts in small volumes of bulk liquid incorporated in the cell design, Fig. 1a and Supplementary Fig. 1. The frequency shifts from the cavity and the bulk markers are of opposite sign, Fig. 1b, c. The superfluid transition within the cavity is sharp, due to the uniformity of cavity height, relative to that achieved in stacked multiple films with a broad distribution of thickness³⁶. The transition temperature depends on whether the surface boundary layer is solid ⁴He or superfluid ⁴He, which determines surface specularity, Fig. 1d. Suppression of the gap by confinement is inferred from the magnitude of the cavity signal frequency shift.

The relatively strong confinement in the 192 nm cavity stabilizes the A phase at all temperatures and pressures, consistent with the phase diagram determined in previous work^{24,36,42}. The orbital angular momentum of the pairs, which defines the orientation of point nodes of the gap in momentum space, lies normal to the cavity surface $\hat{l} = \pm \hat{z}$. The order parameter $\Delta(\mathbf{p}) = \Delta_A(z)(\hat{p}_x + i\hat{p}_y)[|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle]$ ⁴², where \mathbf{p} is the Fermi surface momentum, and z is the position across the cavity. Here Δ_A is the A-phase gap maximum at the Fermi surface equator. In general, the gap has a spatial dependence across the cavity, $\Delta_A(z)$. The static magnetic field $\mathbf{H}_0 = H_0 \hat{z}$ orients the spins along \hat{z} , via the anisotropic magnetic susceptibility. With this relative orientation of spin and orbital angular momentum, the dipolar energy is maximized, accounting for the negative frequency shift observed from the cavity (Fig. 1 and Supplementary Note 1).

Confinement with ⁴He surface plating. We first report measurements in which the sample walls and heat exchanger surfaces were plated with sufficient ⁴He to displace the magnetic solid ³He surface boundary layer, which arises naturally in pure ³He samples⁴³. The plating procedure, which uses a ⁴He surface coverage of 32 $\mu\text{mol m}^{-2}$, results in a non-magnetic localized solid ⁴He surface boundary layer (‘Methods’). In this case, the observed T_c suppression is close to that predicted for purely diffuse scattering. Details of the treatment of surface scattering used in our quasiclassical computations are given in Supplementary Note 3. The results are best fit with specularity $S=0.1$, referred to here as ‘diffuse’, Fig. 2a. The increase in T_c suppression with decreasing pressure arises naturally from stronger effective confinement⁴², Fig. 2b. Subsequently, a

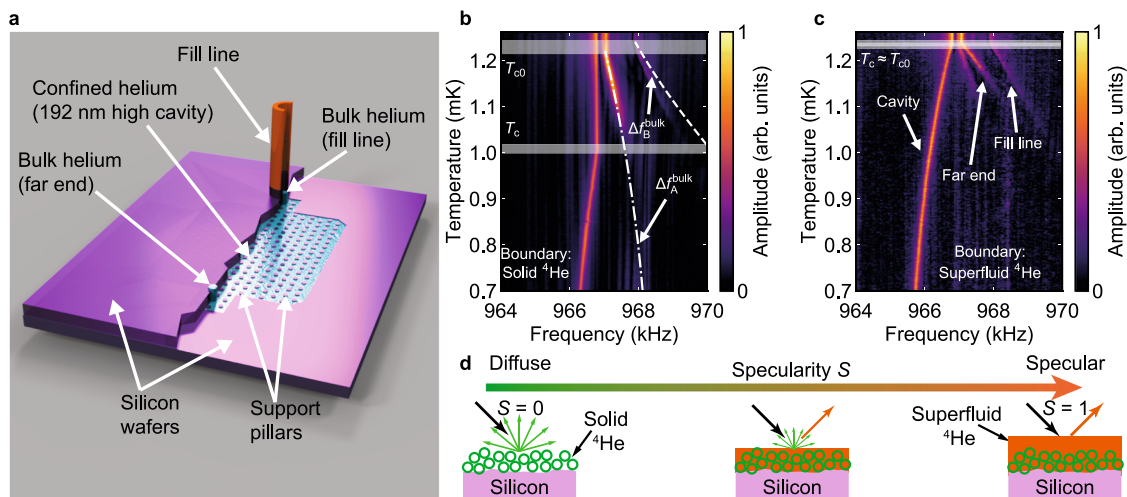


Fig. 1 Experimental cell confining ^3He . **a** Nanofabricated sample cell, cut away to show cavity in lower silicon wafer, bonded to upper wafer. The support posts shown maintain cavity height D under different liquid pressures. The cavity is filled through a fill line via a sintered heat exchanger and cooled through the column of ^3He within it. Small volumes of bulk liquid at each end of the cavity provide markers for the bulk superfluid transition T_{c0} , and eliminate thermometry errors due to temperature gradients. The NMR coil set around the sample is shown in Supplementary Fig. 1. Suitable small magnetic field gradients are used to resolve the NMR response of different regions of the cell, see ‘Methods’. **b, c** NMR signatures of superfluid transition in cavity and bulk markers, for two different surface boundary conditions, at ^3He pressure of 2.46 bar. ^3He -A in the cavity shows a negative frequency shift whereas the bulk markers show positive frequency shift; Δf_A^{bulk} and Δf_B^{bulk} refer to the calculated bulk superfluid frequency shifts of ^3He -A and ^3He -B, respectively, see Supplementary Note 2. The T_c suppression observed with a surface boundary layer of solid ^4He is eliminated by the addition of ^4He to create a superfluid ^4He film at the surface. The white horizontal bands show the measured T_c and T_{c0} including the uncertainties in temperature determination, see ‘Methods’. **d** Schematic illustration of the tuning of surface scattering conditions, parametrised by specular coefficient S , by surface plating atomically smooth silicon with a ^4He film. The green circles represent ^4He atoms, the orange is liquid ^4He , and the arrows indicate the flux of incoming and outgoing ^3He quasiparticles.

thicker ^4He film was formed on the cavity walls to create a surface superfluid film of ^4He (‘Methods’). In this case, we observe an almost complete elimination of T_c suppression, demonstrating close to fully specular scattering, referred to here as ‘specular’, Fig. 2a, b.

In general, the measured frequency shift is related to the spatial average of the suppressed gap within the cavity via $|f^2 - f_L^2| = \zeta \langle \Delta_A^2(z) \rangle$, where ζ is an intrinsic material parameter which is pressure dependent but temperature independent (Supplementary Note 1). In the Ginzburg–Landau regime, sufficiently close to T_{c0} , the A-phase bulk gap maximum Δ_A is given by $\Delta_A^2 = \frac{\Delta C_A}{C_n} (\pi k_B T_{c0})^2 (1 - T/T_{c0})$, where $\Delta C_A/C_n$ is set to the measured heat capacity jump at T_{c0} . This expression thus incorporates strong-coupling corrections to the gap near T_{c0} (Supplementary Note 4). For ‘specular’ boundaries, the measured cavity frequency shift corresponds to the unsuppressed bulk gap, Fig. 2c, and allows determination of the constant ζ at each pressure. For the ‘diffuse’ boundary, using the determined value of ζ , we can precisely infer the gap suppression from the measured frequency shift, independent of uncertainties in material parameters and temperature scale, Supplementary Note 1. We find that the observed gap suppression is also best described by $S=0.1$, establishing the consistency of the experimentally determined gap suppression and T_c suppression within the framework of quasiclassical theory.

Confinement without ^4He surface plating. We now turn to the results where no ^4He preplating was deployed. Rapid exchange between the magnetic surface boundary layer of localized ^3He and the liquid results in a single hybridized NMR line³⁶. The superfluid transition temperature is inferred from analysis of the frequency shift of the hybridized line, which is a weighted average of the internal dipolar frequency shift in the solid ^3He surface boundary layer and that due to superfluidity³⁶ (Supplementary

Note 5). It shows an unexpectedly large T_c suppression, Fig. 3a, significantly exceeding that observed with a solid ^4He boundary layer, and inconsistent with diffuse scattering $S \approx 0$. This result can be phenomenologically described in terms of an effective specularity $S_{\text{eff}} = -0.4$. This approaches the condition for maximal pair-breaking $S = -1$, corresponding to full retroreflection, in which case the phase shift ϕ experienced by the retro-reflected quasiparticle is $\phi = \pi$ for all incoming/outgoing trajectories and surface-bound states accumulate at zero energy, since the excitation energy is given by $E/\Delta = \pm \cos(\phi/2)$ ²⁸, where Δ is the gap along the quasiparticle’s trajectory (see density of states calculation as a function of specularity in Fig. 3c). However, momentum scattering with a preponderance of retroreflection is inconsistent with measurements of boundary slip in viscous transport in the normal state³², and incompatible with the atomically smooth silicon surface.

We invoke magnetic surface scattering to explain this stronger T_c suppression in the presence of the magnetic solid ^3He surface boundary layer. Exchange interaction between quasiparticles and isolated magnetic impurities has been theoretically established to induce additional bound states in superconductors through the Yu–Shiba–Rusinov mechanism⁴⁴. Magnetic scattering by localized ^3He ⁴⁵ strongly influences the observed superfluid phase diagram of ^3He in different aerogels^{46,47}. Here we extend these ideas to consider exchange scattering by the uniform 2D surface layer. We seek processes which generate an excess of zero-energy states over that found for diffuse momentum scattering (Fig. 3b and Supplementary Note 6).

The structure of the order parameter is such that the phase shift ϕ experienced by the scattered quasiparticle, and hence the energy of the surface-bound states, will be influenced by spin-dependent scattering processes. To account for extra pair breaking we need to include quantum spin dynamics of randomly oriented localized quantum spins, allowing for their spin flips. Exchange coupling

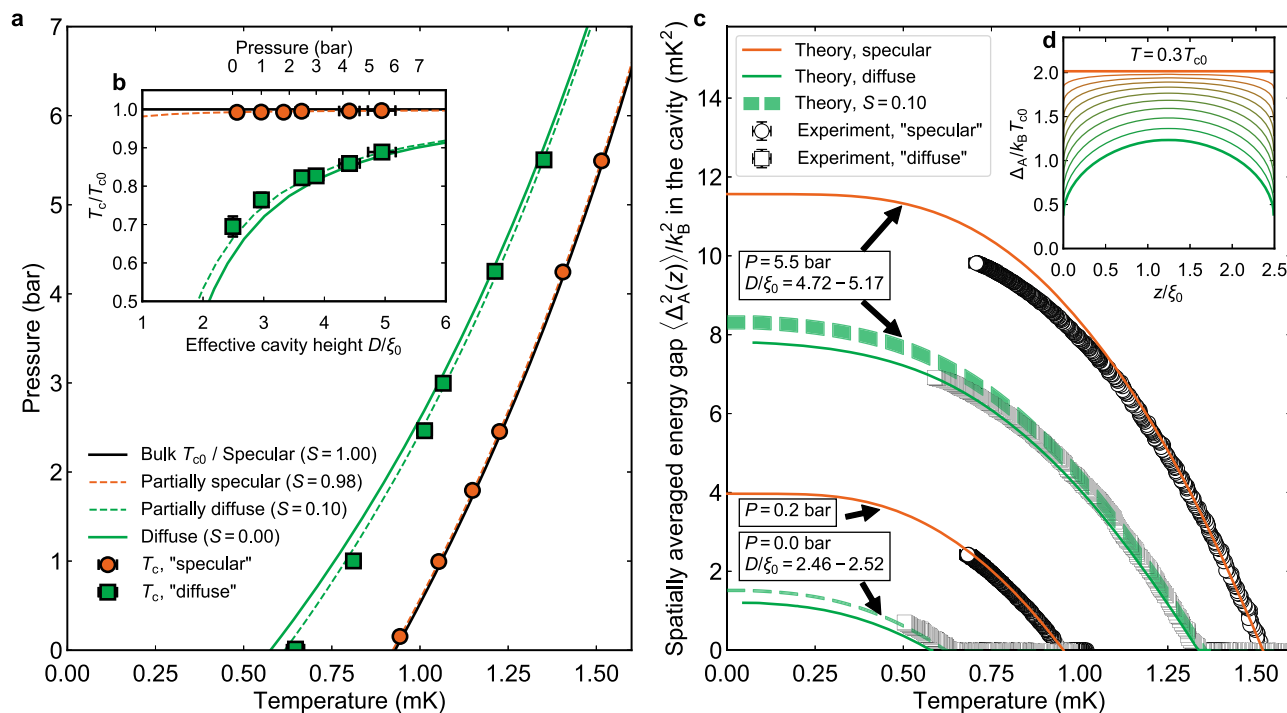


Fig. 2 Suppression of superfluid transition temperature and superfluid gap for different surface scattering conditions. **a** Measured pressure dependence of T_c for close to diffuse (squares) and close to specular (circles) boundary conditions. Full lines show predicted T_c for diffuse (green) and fully specular (black) boundaries, dashed lines are best fits yielding $S = 0.1$ and $S = 0.98$. **b** Suppression of T_c relative to bulk superfluid transition temperature T_{c0} for 'diffuse' boundary steeply increases with confinement. Suppression of T_c for 'specular' boundary is essentially eliminated. The uncertainties in temperature determination and the modelled distortion of cavity height by pressure are used to define the vertical and horizontal error bars, respectively ('Methods'). **c** Spatial average of energy gap $\Delta_A(z)/k_B$ in the cavity, where z is the vertical position in the cavity, inferred from measured frequency shift (Supplementary Note 1), for 'diffuse' (squares) and 'specular' (circles) boundary conditions. All theoretical curves include strong-coupling corrections valid near T_c . The 'diffuse' experiments agree best with theory for $S = 0.1$ (dashed lines, see also Supplementary Fig. 9). The emergent discrepancy between theory and experiment at lower temperatures at 5.5 bar, for both scattering conditions, is in agreement with the expected temperature dependence of strong-coupling corrections to the gap (Supplementary Note 4). The width of the theoretical curves for $S = 0.1$ accounts for errors associated with the weak pressure dependence and uncertainty of cavity height ('Methods'). The theoretical curves for diffuse boundary condition (solid green lines) correspond to the mean value of D/ξ_0 at given pressure, and the theoretical curves for specular boundary condition (solid orange lines) correspond to the bulk energy gaps of $^3\text{He-A}$. **d** The calculated gap profile at zero pressure for specularities between 0 and 1 in intervals of 0.1.

between such spins and incident quasiparticles gives rise to interference between the singlet and triplet scattering channels. In our experimental configuration we find this to be the only viable mechanism resulting in enhanced T_c suppression. We find that for a surface with momentum scattering specularity S , the suppression of T_c corresponds to an effective specularity between bounds $-S \leq S_{\text{eff}} \leq S$, depending on strength of exchange coupling. Thus this process is only detectable for non-diffuse surfaces, but can give rise to T_c suppression exceeding that for a diffuse surface, as observed. To explain the detected $S_{\text{eff}} = -0.4$, we propose that the underlying specularity for momentum scattering from the atomically smooth silicon surface with solid ^3He surface boundary layer should be $S \geq 0.4$.

Discussion

We first focus on a more detailed discussion of the in situ tuning of surface specularity in the context of prior work. We emphasize that the surfaces in our experiment are close to atomically flat ('Methods'), unlike the surfaces used in earlier work. We propose that this extreme smoothness accounts for the partially specular momentum scattering we infer in the absence of ^4He preplating. We suggest that the self-assembled ^3He solid coating at silicon surface does not significantly degrade the smoothness of the liquid–solid boundary.

This contrasts with previous work where diffuse surface scattering ($S = 0$) has been observed in pure ^3He . Diffuse scattering

was inferred from normal state transverse acoustic impedance experiments where the surface roughness of the quartz crystal was of the order of the wavelength of visible light⁴⁰. Suppression of the superfluid T_c of saturated films of pure ^3He was observed using a torsional oscillator, with typical surface roughness of the order of 50 nm⁴⁸. The results agreed with the prediction for fully diffuse surface scattering. This is consistent with our model in which spin-flip magnetic scattering is ineffective when $S = 0$. On the other hand, torsional oscillator studies³³ of surface slip in the normal state with a polished silicon surface of roughness 2 nm found a higher specularity $S \approx 0.2$, confirming the importance of surface roughness.

Our determination of a specularity $S = 0.1$, from superfluid gap and T_c suppression, at the chosen ^4He coverage of $32 \mu\text{mol m}^{-2}$ ('Methods') is in qualitative agreement with prior normal state measurements^{33,35}. These established a rapid increase in specularity for ^4He coverages greater than around 2 layers, attributable to the required coverage for the onset of superfluidity in a ^4He surface film with ^3He overlayer⁴⁹. The absolute values of specularity inferred in the normal state are subject to uncertainty due to imperfect agreement with normal state slip theory³³. We also note the difference between treatment of surface scattering in kinetic transport theory⁵⁰, via phenomenological scattering rates for distribution functions, and quasiclassical theory of the superfluid^{31,41}, via a more microscopic surface scattering matrix. This means that a comparison of boundary scattering parameters is not entirely straightforward.

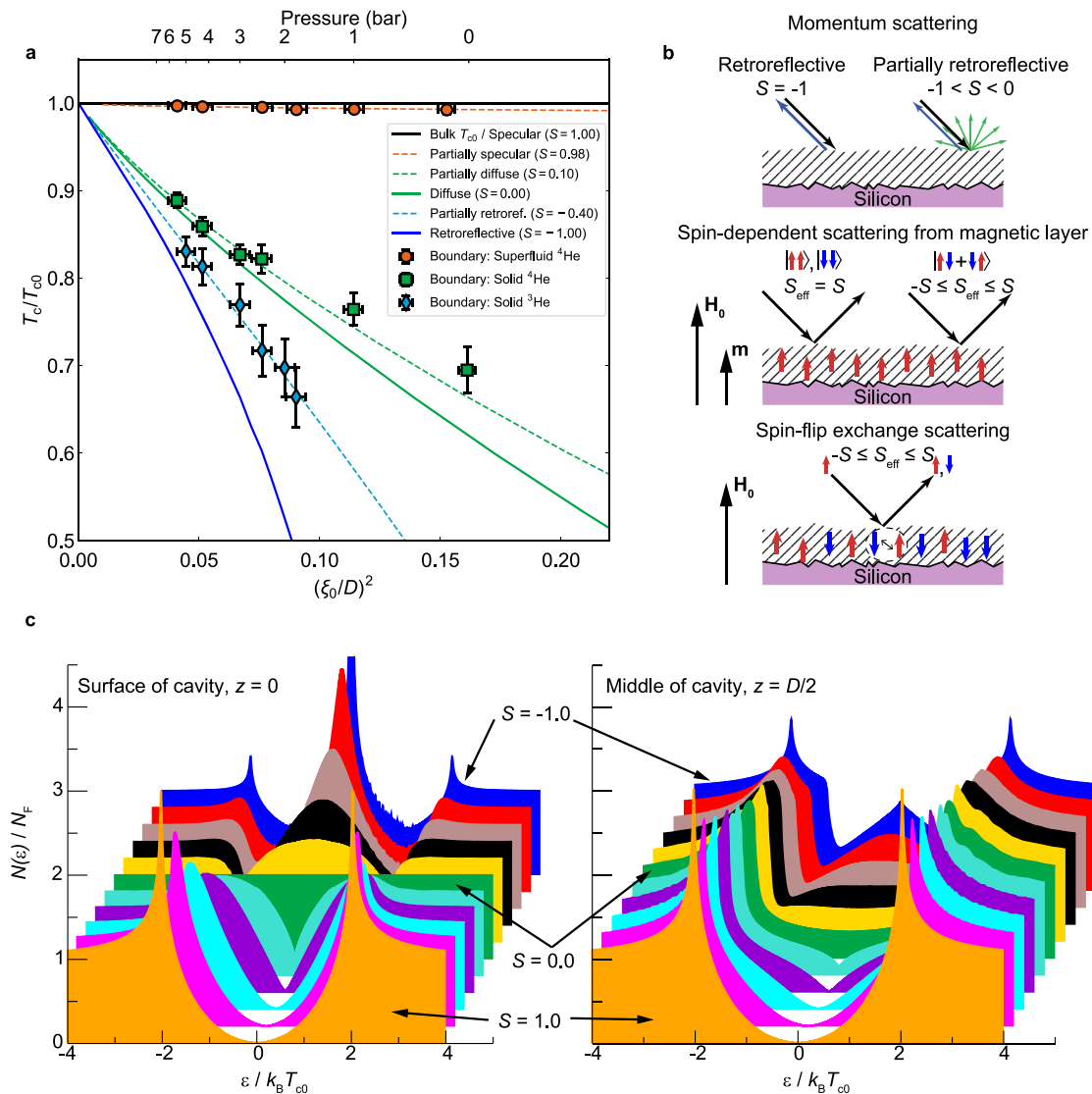


Fig. 3 Increased suppression of superfluid transition temperature in presence of a magnetic solid ^3He surface boundary layer. **a** Suppression of T_c in cavity relative to that of bulk liquid as a function of square of the inverse effective cavity height D/ξ_0 , for solid ^3He boundary (diamonds), solid ^4He boundary (squares), and superfluid ^4He boundary (circles). Full lines show: maximal pair-breaking retroreflection ($S = -1$); diffuse ($S = 0$); fully specular ($S = 1$). Dashed lines show best fits to the data: $S = -0.4, 0.1, 0.98$. For the solid ^3He boundary, T_c is identified from onset of superfluid frequency shift after correcting for background frequency shift arising from magnetic solid layer (Supplementary Note 5). The vertical and horizontal error bars reflect the uncertainties in temperature determination and the height range of the pressure-distorted cavity, respectively ('Methods'). **b** Three candidate scattering mechanisms for negative effective specularity (see also Supplementary Note 6): retroreflection (ruled out by normal state measurements); spin-dependent pair breaking on scattering from a magnetically polarized layer (absent for the relative orientation of surface layer spin polarization \mathbf{m} , surface normal, and spin orientation of A-phase pairs in our set-up); spin-flip exchange scattering (spin polarization of surface layer can be zero). Here the effective specularity S_{eff} is a parameter characterising combined magnetic and momentum scattering, bounded by the specularity S that would arise from momentum scattering alone. **c** Calculated density of states (DOS) $N(\epsilon)$, where ϵ is the quasiparticle energy, of $^3\text{He-A}$ at the surface and in the middle of the cavity, scaled with the normal state DOS, N_f , and averaged over the full Fermi surface. These calculations correspond to pure momentum scattering between specularities $S = -1.0$ and $S = 1.0$, with step size 0.2 between successive colours. The calculations were performed for cavity height $D/\xi_0 = 5$ and temperature $T = 0.2T_{c0}$. The density of zero-energy bound states increases throughout the cavity as the specularity decreases.

The fact that momentum scattering with the pre-plated solid ^4He surface boundary layer is close to diffuse contrasts with the significant specularly we infer in the pure ^3He case from our model of magnetic surface scattering. The replacement of surface ^3He by ^4He relies on the differential surface binding energy of the two isotopes due to differences in zero-point energy. It appears that a ^4He coverage of $32 \mu\text{mol m}^{-2}$, sufficient to displace the ^3He magnetic surface boundary layer but below that required for surface ^4He superfluidity ('Methods'), results in a greater surface roughness than the self-assembled solid ^3He boundary layer. We

suggest that this is attributable to heterogeneity of the surface binding potential across the geometrically smooth silicon surface, which in the presence of the liquid ^3He overburden gives rise to a non-uniform coating of ^4He . At sufficiently high ^4He coverages we establish a superfluid ^4He film covering the entire surface, which creates an ideally smooth equipotential surface and naturally leads to fully specular scattering conditions.

It is clear that for a given set-up the surface scattering can be tuned in situ in a controlled way. Nevertheless, these results provide motivation for further systematic study of the specularly,

over a fine grid of ^4He coverages of the surface boundary layer, to investigate the interplay of momentum and magnetic scattering. This would include detailed measurements in the vicinity of the coverage required to fully displace the ^3He surface boundary layer, and the somewhat higher threshold coverage for the onset of superfluidity in the ^4He surface boundary layer³⁶. Ideally, measurements of superfluid T_c and gap suppression would be coupled with measurements of surface slip in the normal state, and with the development of an independent measurement of surface specularly in the superfluid.

The elimination of gap suppression by specular surfaces we have demonstrated opens the investigation of cavities of arbitrarily small height towards $D \ll \xi_0$ entering the quasi-2D limit, in which thermal and spin analogues of the Quantum Hall effect are predicted^{51,52}. On the other hand, the superfluidity will be completely suppressed in cavities thinner than 100 nm at zero pressure for diffuse scattering, Fig. 2b. Magnetic ^3He boundaries may stabilize new order parameters under confinement and influence surface spin currents. Precise determination of the gap in this case requires measurements in lower magnetic fields, in order to suppress the solid dipolar shift and increase the superfluid frequency shift. This should be possible using broadband SQUID NMR^{53,54}.

Future topological superfluid ^3He mesoscopic devices should provide a new platform for the study of MZMs at well-defined interfaces. The sculpture of the superfluid by confinement will allow the fabrication of hybrid devices based on different ^3He ‘materials’, with clean transparent interfaces, Fig. 4. Design of these platforms is supported by quasiclassical theory, which self-consistently describes the spectrum of surface states, gap suppression, and T_c suppression.

In conclusion, our results show experimentally the sensitivity of the superfluid gap suppression and hence ABS to the details of quasiparticle scattering from the surface. A superfluid ^4He surface boundary layer results in $S = 0.98$, close to fully specular scattering, which eliminates surface ABS in chiral superfluid $^3\text{He-A}$. Whether the specularly can be increased further by a thicker ^4He film, or whether this is a limit determined by solubility of ^3He in the ^4He film, should be resolved by future work. A surface boundary layer of solid ^4He leads to close to diffuse scattering, with a finite density of low energy ABS. These results complement measurements of the transverse acoustic impedance of superfluid $^3\text{He-B}$, which are consistent with the density of surface-bound excitations calculated from quasiclassical theory and their dependence on specularly⁴⁰.

A magnetic solid ^3He surface boundary layer gives rise to a greater suppression of T_c than for diffuse scattering and thus further increases the density of low energy ABS. We show that magnetic scattering from this layer can account for the extra pair breaking. This effect requires a degree of specular momentum scattering and disappears when momentum scattering is diffuse.

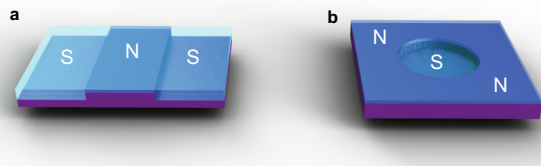


Fig. 4 Topological mesoscopic superfluidity, where confinement tunes ^3He into different material phases, enabling hybrid structures. a An SNS junction, where spatial modulation of cavity height defines SN interfaces. **b** Circular region of higher cavity height defines an isolated mesa of superfluid, cooled through normal liquid in a more confined region. Purple plates represent the silicon wafer, darker blue regions normal liquid ^3He , and lighter more transparent regions superfluid ^3He .

We propose that the extremely smooth walls of our nanofluidic geometry are instrumental in this observation. The surface states of $^3\text{He-A}$ are not topologically protected, unlike its edge states⁵⁵ or the Majorana surface states of $^3\text{He-B}$ ⁵⁶. The calculation of the influence of spin-flip magnetic boundary scattering on the density of states of surface excitations in these two topological superfluids remains an open problem.

In more complex topological superconductor architectures designed to realize, detect, and manipulate MZMs for topological quantum computation, it is necessary to eliminate the excess low energy ABS arising from the interface scattering processes such as we have identified here. Meanwhile, our result is a crucial step in the quest to identify and harness Majoranas in liquid ^3He , the as yet only firmly established topological superconductor/superfluid. More generally, the influence of magnetic degrees of freedom in topological materials⁵⁷ and spin-active interfaces in hybrid superconducting-ferromagnetic spintronics⁵⁸ are both areas of current interest to which the understanding of spin-dependent surface scattering in a spin-triplet superfluid will contribute.

Methods

Silicon nanofluidic cavity fabrication. The experimental cell was fabricated by direct wafer bonding of two silicon wafers. The confinement region and supporting pillars are defined lithographically on one of the wafers using a process similar to that used in a previous generation of cells⁵⁹. The typical rms surface roughness of the silicon surface is 0.1 nm⁶⁰. This is significantly smoother than the mechanically polished silicon surfaces for which surface specularly has been characterized by normal state studies of slip in viscous transport³², potentially promoting specularly of surface scattering even in the absence of a superfluid ^4He film. Deep reactive ion etch (DRIE) is used to create two 300-micron-diameter holes on either side of the confinement region. One acts as a fill line and the other provides a region of bulk helium on the far side of the slab-shaped cavity, Fig. 1a. DRIE is also used to pattern the backside of the wafer to improve the joint between the cell and an external fill line⁶⁰. After all the features are patterned onto the wafers, they are cleaned using a combination of a two-step RCA clean at 75°C followed by immersion in concentrated (49%) HF to remove any oxide or contaminants. The clean patterned wafer is brought into contact with a blank silicon wafer within a wafer aligner, forming a bond between the cavity wafer and the lid. The bond strength is increased and made permanent by an annealing step at 1000°C for 2 h. Successful bonding is confirmed using infra-red imaging and scanning acoustic microscopy. The bonded wafer is diced into individual cells using a diamond saw. A 500-nm-thick silver film is evaporated onto the outside of the bonded wafers to thermalize the cell to the nuclear stage. In order to minimise the effects of differential thermal contraction between the metallic fill line/far-end bulk marker plug and the silicon cell, laser-machined silicon washers are attached around both of the DRIE holes with epoxy (Stycast 1266 mixed with silicon powder), Supplementary Fig. 1. The height of the cavity used in this work was determined to be 192 nm by a profilometer scan on the unbonded wafer. The error in the cavity thickness ± 2 nm was estimated from the distribution in height measured in this way across all the cavities on the unbonded wafer. The maximal distortion of cavity height by pressure is determined by finite element method simulations to be 2.6 nm bar⁻¹, Supplementary Fig. 2. The horizontal error bars in Figs. 2b and 3a reflect the range in cavity height.

NMR measurements. The cooling of the ^3He within the cell, the thermometry, and the SQUID NMR spectrometer were as used in previous work^{24,25}, Supplementary Fig. 1. The helium is cooled via the column of ^3He in the fill line which connects the cell to a sintered silver heat exchanger mounted on a silver plate, connected via a silver rod to the copper nuclear demagnetization stage. A platinum NMR thermometer is mounted on the silver plate. Measurements were made at a ^3He Larmor frequency of 967 kHz, with the static field of around 30 mT applied along the cavity surface normal (\hat{z}). Field gradients were applied to both separate the bulk marker signals from the cavity signal (along \hat{z}) and to resolve the signals from the two bulk markers (along \hat{x} , \hat{y}). The free induction decay following small angle ($3\text{--}10^\circ$) tipping pulses, applied at 10 s intervals, was averaged typically 30 times, and Fourier transformed. To measure the superfluid transitions in the cavity and the bulk markers, the silver plate temperature was swept across the relevant region at a rate below 50 $\mu\text{K h}^{-1}$, Supplementary Fig. 6. The reported values of T_c are averages of several such sweeps. Measurements with tipping pulses of different amplitude enabled a correction to be made for temperature gradients between the ^3He in the cell and the platinum thermometer (Supplementary Note 7). This correction depended on the surface boundary condition, which influenced the boundary resistance of the silver heat exchanger. The uncertainties in the correction introduce an additional temperature error, which has been included in the error bars of all reported temperature values. The temperature gradient across the

cavity is small, and dependent on surface boundary condition; it is determined from the difference between the measured superfluid transition temperatures in the two bulk marker volumes. For the solid ^4He and ^3He surface boundary layers, the difference is around 20 μK , while for the superfluid ^4He surface boundary layer it is at most 2 μK , Supplementary Fig. 6. This gradient is taken into account in determining the error in superfluid transition temperature.

In situ tuning of surface scattering. (i) Magnetic scattering. Pure ^3He , with ^4He impurity concentration <100 ppm, is used to fill the empty cell and silver heat exchanger (surface area 8.03 ± 0.04 m 2 determined by N_2 BET isotherm). This results in a magnetic surface boundary layer of solid ^3He ⁴³. As discussed, the results suggest that this self-assembled solid ^3He layer on the extremely smooth silicon surface gives rise to partially specular momentum scattering which combines with spin-flip magnetic scattering. (ii) Diffuse non-magnetic scattering. In order to displace the naturally occurring magnetic surface boundary layer of ^3He ⁴³, 32 $\mu\text{mol m}^{-2}$ of ^4He was added to the empty cell and silver heat exchanger at 30 K, followed by cooling to below 1 K over 30 h, and a subsequent anneal at 2 K for several hours. This coverage is below that established to be necessary to see a superfluid transition in the surface ^4He layer on a mylar substrate, in the presence of an overburden of ^3He at saturated vapour pressure⁴⁹, and was motivated by previous work³⁶. The sample is cooled to 100 mK before adding ^3He . Under these conditions, the ^3He surface magnetism seen in pure ^3He samples is eliminated, as confirmed by the absence of an observable temperature-dependent susceptibility down to the lowest temperatures achieved. We find that the remaining momentum scattering is close to diffuse. (iii) Specular scattering. To create specular scattering conditions from the previous ^4He surface plating conditions, the cell is pumped at 1.5 K, leaving a residual solid ^4He 'layer' on the surfaces. Then more ^4He is added into the cell/heat exchanger. Subsequently, the helium pumped out in the previous step is restored. The sample is slowly cooled into the mK range, during which all the ^4He forms a surface film of solid ^4He with a superfluid ^4He overlayer. With nominal surface ^4He coverage in the range 68–139 $\mu\text{mol m}^{-2}$, we always detect the same specularly, consistent with previous work³², which found evidence for surface scattering close to specular for surface film coverages in excess of 60 $\mu\text{mol m}^{-2}$. The quoted surface coverages assume that the thin ^4He films coat the heat exchanger and sample volume uniformly. We note that differences in heat exchanger structures and surface area determination, as well as varying methodologies for adding ^4He into the system may lead to systematic differences between nominal sample surface ^4He coverages between different set-ups in different laboratories.

Theoretical calculation of gap suppression. Theoretical calculation of gap suppression is made using quasiclassical theory that can systematically include spatial variations of the order parameter and Andreev scattering process. The boundary conditions on the propagator are incorporated using the random S-matrix scattering model (Supplementary Note 3).

Data availability

The T_c and gap suppression data and calculations that support the findings of this study are available in Figshare, <https://doi.org/10.17637/rh.1277762061>.

Received: 20 April 2020; Accepted: 28 January 2021;

Published online: 10 March 2021

References

- Vollhardt, D. & Wölfle, P. *The Superfluid Phases of Helium 3* (Dover, 2013).
- Leggett, A. J. A theoretical description of the new phases of liquid ^3He . *Rev. Mod. Phys.* **47**, 331–414 (1975).
- Volovik, G. E. *The Universe in a Helium Droplet* (Clarendon, 2003).
- Mizushima, T. et al. Symmetry-protected topological superfluids and superconductors—from the basics to ^3He . *J. Phys. Soc. Jpn.* **85**, 74 (2016).
- Ando, Y. & Fu, L. Topological crystalline insulators and topological superconductors: from concepts to materials. *Annu. Rev. Condens. Matter Phys.* **6**, 361–381 (2015).
- Qi, X.-L. & Zhang, S.-C. Topological insulators and superconductors. *Rev. Mod. Phys.* **83**, 1057–1110 (2011).
- Schnyder, A. P., Ryu, S., Furusaki, A. & Ludwig, A. W. W. Classification of topological insulators and superconductors in three spatial dimensions. *Phys. Rev. B* **78**, 195125 (2008).
- Mackenzie, A. P., Scaffidi, T., Hicks, C. W. & Maeno, Y. Even odder after twenty-three years: the superconducting order parameter puzzle of Sr_2RuO_4 . *npj Quantum Mater.* **2**, 40 (2017).
- Joynt, R. & Taillefer, L. The superconducting phases of UPt_3 . *Rev. Mod. Phys.* **74**, 235–294 (2002).
- Sasaki, S. et al. Topological superconductivity in $\text{Cu}_x\text{Bi}_2\text{Se}_3$. *Phys. Rev. Lett.* **107**, 217001 (2011).
- Ran, S. et al. Nearly ferromagnetic spin-triplet superconductivity. *Science* **365**, 684–687 (2019).
- Sato, M. & Ando, Y. Topological superconductors: a review. *Rep. Prog. Phys.* **80**, 076501 (2017).
- Alicea, J. New directions in the pursuit of Majorana fermions in solid state systems. *Rep. Prog. Phys.* **75**, 076501 (2012).
- Sarma, S. D., Freedman, M. & Nayak, C. Majorana zero modes and topological quantum computation. *npj Quantum Inf.* **1**, 15001 (2015).
- Beenakker, C. W. J. Search for Majorana fermions in superconductors. *Annu. Rev. Condens. Matter Phys.* **4**, 113–136 (2013).
- Aasen, D. et al. Milestones toward Majorana-based quantum computing. *Phys. Rev. X* **6**, 031016 (2016).
- Deng, M. T. et al. Majorana bound state in a coupled quantum-dot hybrid-nanowire system. *Science* **354**, 1557–1562 (2016).
- Karzig, T. et al. Scalable designs for quasiparticle-poisoning-protected topological quantum computation with Majorana zero modes. *Phys. Rev. B* **95**, 235305 (2017).
- Lutchyn, R. M. et al. Majorana zero modes in superconductor–semiconductor heterostructures. *Nat. Rev. Mater.* **3**, 52–68 (2018).
- Ren, H. et al. Topological superconductivity in a phase-controlled Josephson junction. *Nature* **569**, 93–98 (2019).
- Fornieri, A. et al. Evidence of topological superconductivity in planar Josephson junctions. *Nature* **569**, 89–92 (2019).
- Liu, C.-X., Sau, J. D., Stanescu, T. D. & Das Sarma, S. Andreev bound states versus Majorana bound states in quantum dot-nanowire-superconductor hybrid structures: trivial versus topological zero-bias conductance peaks. *Phys. Rev. B* **96**, 075161 (2017).
- Setiawan, F., Wu, C.-T. & Levin, K. Full proximity treatment of topological superconductors in Josephson-junction architectures. *Phys. Rev. B* **99**, 174511 (2019).
- Levitin, L. V. et al. Phase diagram of the topological superfluid ^3He confined in a nanoscale slab geometry. *Science* **340**, 841–844 (2013).
- Levitin, L. V. et al. A nuclear magnetic resonance spectrometer for operation around 1 MHz with a sub-10-mK noise temperature, based on a two-stage dc superconducting quantum interference device sensor. *Appl. Phys. Lett.* **91**, 262507 (2007).
- Levitin, L. V. et al. Surface-induced order parameter distortion in superfluid ^3He -B measured by nonlinear NMR. *Phys. Rev. Lett.* **111**, 235304 (2013).
- Levitin, L. V. et al. Evidence for a spatially modulated superfluid phase of ^3He under confinement. *Phys. Rev. Lett.* **122**, 085301 (2019).
- Vorontsov, A. B. Andreev bound states in superconducting films and confined superfluid ^3He . *Philos. Trans. A Math. Phys. Eng. Sci.* **376**, 2015.0144 (2018).
- Serene, J. W. & Rainer, D. The quasiclassical approach to superfluid ^3He . *Phys. Rep.* **101**, 221–311 (1983).
- Halperin, W., Choi, H., P. Davis, J. & Pollanen, J. Impurity effects of aerogel in superfluid ^3He . *J. Phys. Soc. Jpn.* **77**, 111002 (2008).
- Nagato, Y., Yamamoto, M. & Nagai, K. Rough surface effects on the p-wave Fermi superfluids. *J. Low. Temp. Phys.* **110**, 1135–1171 (1998).
- Tholen, S. M. & Parpia, J. M. Slip and the effect of ^4He at the ^3He -silicon interface. *Phys. Rev. Lett.* **67**, 334–337 (1991).
- Tholen, S. M. & Parpia, J. M. Effect of ^4He on the surface scattering of ^3He . *Phys. Rev. B Condens. Matter* **47**, 319–329 (1993).
- Wada, Y. et al. Broadening of the surface Andreev bound states band of superfluid ^3He -B on a partially specular wall. *Phys. Rev. B* **78**, 214516 (2008).
- Murakawa, S. et al. Strong suppression of the Kosterlitz-Thouless transition in a ^4He film under high pressure. *Phys. Rev. Lett.* **108**, 025302 (2012).
- Freeman, M. R. & Richardson, R. C. Size effects in superfluid ^3He films. *Phys. Rev. B* **41**, 11011–11028 (1990).
- Kim, D. et al. Boundary condition on superfluid ^3He as altered by ^4He interfacial layer. *Phys. Rev. Lett.* **71**, 1581–1584 (1993).
- Murakawa, S. et al. New anomaly in the transverse acoustic impedance of superfluid ^3He -B with a wall coated by several layers of ^4He . *Phys. Rev. Lett.* **103**, 155301 (2009).
- Murakawa, S. et al. Surface Majorana cone of the superfluid ^3He -B phase. *J. Phys. Soc. Jpn.* **80**, 013602 (2011).
- Okuda, Y. & Nomura, R. Surface Andreev bound states of superfluid ^3He and Majorana fermions. *J. Phys.-Condens. Matter* **24**, 343201 (2012).
- Nagai, K., Nagato, Y., Yamamoto, M. & Higashitani, S. Surface bound states in superfluid ^3He . *J. Phys. Soc. Jpn.* **77**, 111003 (2008).
- Vorontsov, A. B. & Sauls, J. A. Thermodynamic properties of thin films of superfluid ^3He -A. *Phys. Rev. B* **68**, 064508 (2003).
- Ahonen, A. I. et al. Boundary magnetism in liquid ^3He at very low temperatures. *J. Phys. C: Solid State Phys.* **9**, 1665–1672 (1976).
- Balatsky, A. V., Vekhter, I. & Zhu, J.-X. Impurity-induced states in conventional and unconventional superconductors. *Rev. Mod. Phys.* **78**, 373–433 (2006).

45. Mineev, V. P. Influence of exchange scattering on superfluid ^3He states in nematic aerogel. *Phys. Rev. B* **98**, 014501 (2018).
46. Dmitriev, V. V., Soldatov, A. A. & Yudin, A. N. Effect of magnetic boundary conditions on superfluid ^3He in nematic aerogel. *Phys. Rev. Lett.* **120**, 075301 (2018).
47. Zimmerman, A. M., Nguyen, M. D., Scott, J. W. & Halperin, W. P. Effect of magnetic impurities on superfluid ^3He . *Phys. Rev. Lett.* **124**, 025302 (2020).
48. Xu, J. & Crooker, B. C. Very thin films of ^3He —a new phase. *Phys. Rev. Lett.* **65**, 3005–3008 (1990).
49. McQueeney, D., Agnolet, G. & Reppy, J. D. Surface superfluidity in dilute ^4He - ^3He mixtures. *Phys. Rev. Lett.* **52**, 1325–1328 (1984).
50. Einzel, D., Wölfle, P. & Hirschfeld, P. J. Transverse surface impedance of pair-correlated Fermi liquids. Application to normal and superfluid ^3He . *J. Low. Temp. Phys.* **80**, 31–68 (1990).
51. Volovik, G. E. An analog of the quantum Hall effect in a superfluid ^3He film. *JETP (USSR)* **67**, 1084 (1988).
52. Volovik, G. E. *Exotic Properties Of Superfluid ^3He* (World Scientific, 1992).
53. Arnold, F. et al. Application of low frequency SQUID NMR to the ultra-low temperature study of atomically layered ^3He films adsorbed on graphite. *J. Phys.: Conf. Ser.* **568**, 032020 (2014).
54. Körber, R. et al. Nuclear magnetic resonance on room temperature samples in nanotesla fields using a two-stage dc superconducting quantum interference device sensor. *Appl. Phys. Lett.* **91**, 142501 (2007).
55. Sauls, J. A. Surface states, edge currents, and the angular momentum of chiral p-wave superfluids. *Phys. Rev. B* **84**, 214509 (2011).
56. Wu, H. & Sauls, J. A. Majorana excitations, spin and mass currents on the surface of topological superfluid $^3\text{He-B}$. *Phys. Rev. B* **88**, 184506 (2013).
57. Li, H. et al. Dirac surface states in intrinsic magnetic topological insulators EuSn_2As_2 and $\text{MnBi}_{2n}\text{Te}_{3n+1}$. *Phys. Rev. X* **9**, 041039 (2019).
58. Eschrig, M. Theory of Andreev bound states in S-F-S junctions and S-F proximity devices. *Philos. Trans. A Math. Phys. Eng. Sci.* **376**, 2015.0149 (2018).
59. Dimov, S. et al. Anodically bonded submicron microfluidic chambers. *Rev. Sci. Instrum.* **81**, 013907 (2010).
60. Zhelev, N. et al. Fabrication of microfluidic cavities using Si-to-glass anodic bonding. *Rev. Sci. Instrum.* **89**, 073902 (2018).
61. Heikkinen, P. J. et al. Accompanying data for “Fragility of surface states in topological superfluid ^3He ”. *figshare* <https://doi.org/10.17637/rh.12777620> (2020).

Acknowledgements

We thank M. Eschrig, J.A. Sauls, and S. Simon for helpful discussions. This work was supported by EPSRC grants EP/J022004/1 and EP/R04533X/1, NSF grants DMR-1708341 and DMR-2002692, and the European Union’s Horizon 2020 Research and Innovation Programme, under Grant Agreement no. 824109 (European Microkelvin Platform). Fabrication was carried out at the Cornell Nanofabrication Facility (CNF) with assistance and advice from technical staff. CNF is supported in part by the NSF

through ECCS-1542081. Measurements we made at the London Low Temperature Laboratory, supported by technical staff, in particular Richard Elsom, Ian Higgs, Paul Bamford, and Harpal Sandhu.

Author contributions

Experimental work was carried out by P.J.H. and L.V.L. with contributions from A.C. The nanofabricated cells were prepared and assembled by N.Z., X.R. and A.C. X.R. carried out the FEM simulations of the cell. The analysis and presentation was carried out by P.J.H., L.V.L. and A.V. with contributions from J.S. A.V. performed calculations of gap profile and superfluid transition temperature. A.V. developed the theory of magnetic scattering with contributions from P.S. The work at Cornell was supervised by J.M.P., and the work in London was supervised by J.S., A.C. and L.V.L., who had the leading roles in formulating the research. P.J.H., J.S. and A.V. had leading roles in writing the paper, with contributions from all authors.

Competing interests

The authors declare no competing interests.

Additional information

Supplementary information The online version contains supplementary material available at <https://doi.org/10.1038/s41467-021-21831-y>.

Correspondence and requests for materials should be addressed to P.J.H. or J.S.

Peer review information *Nature Communications* thanks the anonymous reviewer(s) for their contribution to the peer review of this work. Peer reviewer reports are available.

Reprints and permission information is available at <http://www.nature.com/reprints>

Publisher’s note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article’s Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article’s Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>.

© The Author(s) 2021